

Characterizing landslide dynamics from time-lapse time domain induced polarization and ground based imaging, case study of the MontGombert landslide (French, Alps)

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Abstract

10 Electrical Resistivity Tomography (ERT) is a common tool to study landslides morphology and dynamics. Landslides dynamics is strongly related to clay content, pore pressure and permeability distribution. Standard Direct Current (DC) resistivity method often fails at providing information about permeabilities because materials with low resistivities can both exhibits high or low permeability especially according to their clay content and the relative importance
15 of surface conduction compared to electrolytic conduction. Induced polarization (IP) have proven to be a useful complementary tool to assess petrophysical properties of the rocks. Up to now, it has been rarely used in landslides studies and not yet in landslides monitoring. This study aims at demonstrating the usefulness and added value of induced polarization as a complement to traditional
20 ERT measurements in the case of landslide characterization and monitoring. We realized 3 time domain IP (TDIP) surveys at three different time steps on an active landslide in the French Alps. A time-lapse inversion of the resistivity and chargeability values is realized. The spectral IP parameters are extracted
25 through a Debye Decomposition. The results are compared to surface displacements obtained by processing optical images. We demonstrated that the use

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of TDIP is of crucial importance to estimate the extend of unstable volumes within the landslide. The results are in very good agreement with the data provided by the piezometers and inclinometers. The use of a general law to link IP variables and permeability enables to estimate the spatial distribution of the permeability within the landslide and its evolution through the seasons.

Keywords:

Landslides, electrical tomography, induced polarization, permeability, surface displacement, risk assessment

1. Introduction

Nowadays, while climate is changing, society and land-use are evolving and extending, landslides occurrences are increasing. Understanding landslide processes and the risk associated is becoming a major economic and political stake, especially in mountainous areas (Malet et al., 2007; Marescot et al., 2008; Gance et al., 2016; Lajaunie et al., 2019; Scoppettuolo et al., 2020; Desrues et al., 2022). Landslide deformation monitoring is one type of approach to evaluate slope stability and set up early warning techniques. It implies surface deformation monitoring, satellite or image based approaches and/or deformation modelisation approaches (Stumpf et al., 2017; Lacroix et al., 2018; Desrues et al., 2019, 2022; Cascini et al., 2022).

Full landslides characterization consists in evaluating the potential unstable volumes implicated and defining the triggering processes. Changes in water content of the subsurface material, the associated pore water pressure and matrix suction variations are the main involved triggering parameters underlying the strong influence of heavy rainfalls and sudden snow melts in landslides dynamics (Malet et al., 2005; Eberhardt et al., 2005; Guzzetti et al., 2008; Lu et al., 2010; Ponziani et al., 2012; Yao et al., 2019). Geophysical approaches, and especially electrical methods due to their sensitivity to water content, have proven to be

efficient for structure characterization, soil moisture content characterization
55 and evolution through time (Marescot et al., 2008; Hibert et al., 2012; Wilkin-
son et al., 2016; Gance et al., 2016; Hellman et al., 2017; Uhlemann et al., 2017;
Lajaunie et al., 2019; Revil et al., 2020; Holmes et al., 2020). If moisture content
is a main factor controlling landslide stability assessment, the distribution of clay-
ish areas acting as impermeable barriers is also a key to understand landslides
60 dynamics. Recently Uhlemann et al. (2017); Chambers et al. (2014) proposed
a moisture content evaluation based on the measure of the electrical resistivity
and the consideration of surface conduction based on laboratory calibration of
a Waxman and Smith type relationship.

Considering the importance of clay content estimation and distribution in
65 understanding landslides, it turns out logical to introduce induced polarization
(IP) measurements to classical geoelectrical surveys. IP measurements can be
performed both in frequency (FDIP) and in time (TDIP) domains. They can
both lead to information about the amplitude of the polarization and the re-
laxation time. The IP spectral content hold information about petrophysical
70 and hydrological characteristics such as the hydraulic conductivity (Binley and
Kemna, 2005; Tong et al., 2006a; Titov et al., 2010; Kemna et al., 2012) or the
clay content (Slater and Lesmes, 2002; Slater and Glaser, 2003). Up to now,
very few studies present the results of field induced polarization on landslides
(Marescot et al., 2008; Revil et al., 2020; Orozco et al., 2022). Revil et al. (2020)
75 applied petrophysical models in order to estimate water content, cation exchange
capacity and permeability from laboratory calibration to a clayey landslide. Sim-
ilarly to this study, most of the field IP and almost all the field IP realized on
landslides consider only the chargeability parameter depriving themselves of the
spectral information. Recently, Orozco et al. (2022) has performed both TDIP
80 and FDIP on a landslide and has evidenced that hydraulic properties estimated

from spectral IP content where of fundamental importance for landslide characterization. Information about the spectral content of the IP phenomenon are usually retrieved from FDIP measurements due to their higher accuracy and larger acquisition bandwidth. Such acquisitions besides being time greedy, have
85 a lower penetration depth on the field (Maurya et al., 2018; Martin et al., 2020). In the past 10 years, the use of TDIP to retrieve IP spectral information has been developing. Generally, the spectral properties are extracted via a Cole Cole, a Pelton or a maximum phase angle model inversion (Gazoty et al., 2012; Fiandaca et al., 2013; Doetsch et al., 2015; Fiandaca et al., 2018; Lévy et al.,
90 2019a; Martin et al., 2020; Orozco et al., 2022). The latter imply some hypothesis as for the shape of the phase spectra. To avoid such constraints, TDIP decays can also be fitted via a Debye Decomposition (DD) (Tarasov and Titov, 2007; Nordsiek and Weller, 2008; Martin et al., 2020, 2021). To our knowledge the DD approach has never been applied on landslides, and the spectral content
95 evolution through time in landslides has not yet been investigated.

In this paper, we propose a time-lapse ERT and IP monitoring of a landslide situated in the French Alps. The full decay of the IP curve will be considered in order to get information about the relaxation times via a Debye Decomposition.
100 We propose to associate the results to the surface deformation obtained by photogrammetric monitoring. We show that the IP information is important as an added value for landslides geometry characterization, time-lapse monitoring, dynamic processes understanding and unstable volumes quantification.

2. Geological settings and major stakes

105 The Montgombert landslide (MGL) is located in the French part of the alps in the gorges of the Arly river (Savoie, France) near the city of Ugine (Figure 1).

The landslide is perched 100 m above the road RD1212. This road is an essential trunk road as it deserves some major ski resorts in the area and link cities in the pays du Mont Blanc to bigger cities. The part of the RD1212 crossing the gorges of the Arly has been regularly exposed to landslides, rock falls and flooding generating regular traffic interruptions and road works. All the Eastern part of the gorges de l'Arly are prone to landslides and are composed of hercynian metamorphic micaschists from the Serie Satinée. The MGL originates in the generalized toppling failure of the micaschists and their transformation into clayey micaschists rock slides. Further on the East, the micaschists are covered by mesozoic rocks and glacial sediments on the Plateau des Saisies. The latter forms significant local aquifers that generate permanent water circulation into the massif. Water is drained by the surface hydrological network and also diffusively through the rock slides. The MGL has been particularly active since the spring 2016 and strongly monitored ever since. Piezometers and inclinometers have been installed. Surface displacements are monitored thanks to automatic theodolite and optical image monitoring installations (Desrues et al., 2019). The slope is the order of 35° . Geophysical studies have been performed in 2016 and 2017. The most active part of the landslide is about 100 m width and 200 m long and its depth is evaluated to 15 to 30 m. The piezometric level at the top of the active part of the landslide is 18 to 16 m below the topographic surface in spring and 30 to 31 m at the end of the summer. The surface displacements show that movements are more important in the lower part of the landslide and can reach 25 m within a year. The inclinometer exhibits 3 main sliding surfaces at 8, 15 and 30 m depth, the latter being consistent with the water level observed in the piezometer. The landslide is water sensitive and is particularly active in the spring following heavy rains and/or fast snow melting. The particularly rainy winter 2017-2018 generated a sudden reactivation of the MGL.

Open fractures up to 2 m wide and 10 m long have been produced, 7000 m³ of
135 rock debris have reached the road and several hundreds of thousands cubic me-
ters of material have moved from 50 cm to 15 m at the surface of the landslide
and stayed in equilibrium on the slope. As a consequence, the administration of
the French county of Savoie decided to finance more geological, hydrogeological
and geophysical studies to better understand the structure of the landslide and
140 the mechanisms driving its instabilities.

3. Data acquisition and pre-processing

3.1. Photogrammetric data

A system of topographic monitoring is in place on the studied site (Figure
1). We placed 26 topographic targets in 2016 and surveyed their positions with
145 an automatic theodolithe. Following technical issues and some targets loss, the
installation had to be reset and has been operational from February 2021 to
November 2021 and further removed. We installed 2 digital single-lens reflex
cameras (Canon 100D) of 18 MPix connected to a Paratronic LNS datalogger
and separated of 362 m on the northern opposite side (right bank) of the valley
150 enabling a full view of the most active part of the landslide. The cameras
have been installed in May 2019 and are operating over since. The distance
between the camera and the landslide is 550 m, the equivalent size of one pixel
is 5 cm and the theoretical displacement detection limit is 0.5 cm. In order
to avoid problems due to shades, 3 pictures per day are taken at 11:00, 12:00
155 and 13:00 (UTM). The calibration of the cameras has been performed thanks
to 10 georeferenced targets placed on the landslide and a high resolution Lidar
obtained at the moment of the installation of the cameras. We used the images
from the 18th of April 2021, the 16th of September 2021 and the 12th of April
2022. Each pair of images has been processed.

160 *3.2. Geoelectrical data*

We acquired 4 electrical ERT and IP profiles in the active part of the MGL. A 5 m spacing multiple gradient acquisition protocol with a double cable configuration (2×32 stainless steel electrodes) have been performed in order to maximize signal/noise ratio and prevent electromagnetic coupling effects within the cables (further details about this acquisition procedure can be found in Dahlin et al. 165 (2002); Lévy et al. (2019a)). We realized 3 campaigns in April 2021, September 2021 and May 2022. The studied area being particularly steep and unstable, the electrodes have not been left in place from one campaign to the other. At each campaign, the precise position of each electrode is taken with a real time 170 kinematic GPS (Leica GS07 high precision GNSS with a RTK connection) with a 5 cm accuracy. In April 2021, resistivity and induced polarization data were acquired using an ABEM Terrameter SAS 4000. For IP data, 6 time windows (100, 200, 400, 800, 1600, 3200 ms) were set with a time delay of 50 ms to make sure no electromagnetic coupling is recorded on the early times, on-time and 175 off-time were equal and set to 8 s. To reduce the acquisition time in September 2021 and May 2022, a Terrameter LS2 have been used in 100% duty cycle (Olsson et al., 2015). The fullwaveform data enabled a first data quality check, manual inspection and removal of noisy data (Martin et al., 2020). Injection time was set to 8 s and 13 time windows have been recorded with a 50 ms 180 time delay (100, 100, 100, 100, 100, 100, 100, 100, 100, 200, 400, 800, 1600, 3200 ms). Unfortunately, the chargeability values are shifted from 50% duty cycle to 100% duty cycle and we can not ensure the comparability of these data and their decays between April 2021 and the further campaigns (Olsson et al., 2015). This is why we will not further use the April 2021 chargeability data for TDIP 185 time-lapse inversion. A total of 1053, 996 and 977 resistivity and chargeability raw data have been obtained in April 2021, September 2021 and May 2022 re-

spectively. We measured the stacking errors - or variation coefficient - in % over at least 4 stacks. The median value and standard deviation value of these errors are 0.72 and 1.55 % in April 2021, 0.89 and 0.17 % in September 2021 and 0.085 and 0.12 % in May 2022. The errors are less spread when using the LS2 which can be explained by a more recent manufacture and a better optimized internal quality. The stacking errors in September are higher than those obtained in May which results from the quality of the contact resistances ($\approx 1000 \text{ Ohm}$ with a maximum of 1700 in September and $\approx 700 \text{ Ohm}$ in May). Negative resistivities, values with stacking errors superior to 2%, 1.2% and 1% for April 2021, September 2021 and May 2022 as well as outliers values are removed from the resistivity data set. Chargeability decays are checked and filtered such as negative chargeabilities at first step, decays which do not fit an exponential decay and non-monotonic decreasing curves are removed. A total of 1028, 895 and 969 resistivity data and 594, 747 and 905 chargeability decays are kept for further inversion (April 2021, September 2021 and May 2022 respectively).

4. Methods

4.1. Electrical and induced polarization methods

Classical ERT measurements consist in injecting current through a pair of electrodes and measuring the resulting voltage through a pair of receiving electrodes leading to a resistance value. However, the subsurface materials can also act as capacitors, i.e. they can reversibly store electrical charges when an external current is applied. This effect is also known as induced polarization effect or chargeability. Charges accumulation and depletion generate a secondary electrical field. In time domain (TD), this lead to a voltage decay when the current is switched off. The characteristics of the decay curve are fundamentals to define the IP effect not only in terms of magnitude but also in terms of slope

and relaxation time (Gazoty et al., 2012; Martin et al., 2021). The overall conductivity of soils (inverse of resistivity) is the result of three main effects : the matrix conduction, which can be neglected if no metallic particles or electronic
 215 conductors are present, the electrolytic conduction, i.e. the conduction of the charges within the fluid present in the pores, and the surface conduction. i.e. conduction processes happening at the fluid/grain interface. In materials characterized by fine grains like clays, surface conduction can not be neglected. When
 220 no metallic particles are present the conductivity is a complex value that can be expressed as :

$$\sigma^* = [\sigma_{el} + \sigma'_s] + i\sigma''_s \quad (1)$$

with σ_{el} the electrolytic conductivity and σ_s the surface conductivity, both in $S.m^{-1}$. Capacitive properties are related to the imaginary part of the surface conductivity (Kemna et al., 2012). The material used for TDIP measurements is
 225 the same as for standard ERT measurements. Practically, a sequence of pulses of constant current and opposite signs with or without pauses between them (50% versus 100% duty cycles (Olsson et al., 2015; Madsen et al., 2017)) are injected into the ground. Partial apparent chargeabilities are defined as (Slater and Lesmes, 2002; Gazoty et al., 2012; Revil et al., 2020):

$$m_i = \frac{1}{(t_{i+1} - t_i)} \int_{t_i}^{t_{i+1}} \eta(t) dt \quad (2)$$

230 where $\eta(t) = \frac{V_s(t)}{V_{DC}}$ in V/V is the ratio between the secondary potential measured during the off-time and the potential measured between the potential electrodes just before the current is switched off and t_{i+1} and t_i define the beginning and ending of the window over which the signal is integrated. The normalized chargeability is defined as $m_n = \frac{m}{\rho_0}$, where m is the chargeability

235 and ρ_0 the measured DC resistivity. This variable is considered as analogical to
the imaginary part of the conductivity, σ'' , in the FD (Slater and Lesmes, 2002;
Gurin et al., 2013) and can be considered as a proxy of surface conductivity in
sedimentary rocks (Slater and Lesmes, 2002; Revil et al., 2017).

4.2. *Geoelectrical inversion including induced polarization data*

240 The inversion of the electrical and TDIP data have been performed thanks
to the open-source package pyGIMLI (Günther et al., 2006; Rücker et al., 2006,
2017). As a first step, the inversion of the resistivity data is performed us-
ing a Gauss-Newton minimization algorithm implemented on an irregular grid
(Günther et al., 2006; Rücker et al., 2006). A gaussian distribution of the error
245 estimate on the data is considered. The errors are defined as the sum of the
measurement accuracy, resulting from the measurement of stacking errors (Tso
et al., 2017) and a fixed percentage of the measure (Günther et al., 2006). The
stopping criteria is based firstly on the χ^2 criteria and then on the decrease of
the cost function. The model resolution is estimated via the coverage approach
250 developped in Günther et al. (2006). Considering the approach of Seigel (1959),
the apparent chargeabilities are then inverted to get a chargeability model (Mar-
tin et al., 2020; Revil et al., 2020). One model per time gate is extracted and
a smoothing constrained is applied along the time axes to ensure the models
decays to be smooth (Martin et al., 2020). Once the inversion of the different
255 time gates is realized, decays can be further fitted.

4.3. *Time-lapse inversion methodology*

To track the evolution and variations of resistivity and chargeability values
and infer the influence of the water content in the MGL, the ERT and IP ex-
periments are also conducted as a time-lapse experiment. The accessibility and
260 the stability of the studied area remain challenging. Some areas at the base of

the studied site became too unstable from one campaign to the other to ensure
 a safe access implying some electrodes relocation. As a result, our method-
 ology for the time-lapse inversion is similar to the methodology proposed by
 Uhlemann et al. (2017). A very precise mesh is build on a 20 *cm* z-precision
 265 Digital Elevation Model (DEM) and includes one node per electrode position
 all campaigns combined. The electrodes positions are updated for each time
 step and a reference model is used to constrain the inversion. The same inver-
 sion mesh is used for all the inversions as no significant topographic changes
 compared to the grid size that may affect the resistivity model are observed
 270 during our time study. Above 0°C, conductivity and normalized chargeability
 increase linearly of about 2% per °C (Hayley et al., 2007; Revil et al., 2017;
 Uhlemann et al., 2017; Holmes et al., 2020). In order to avoid misinterpretation
 about water content variation, the influence of the seasonal temperature needs
 to be removed. Brunet et al. (2010) proposed to model the temperature as a
 275 function of depth, annual mean temperature, yearly amplitude variation of the
 air temperature, a characteristic penetration depth, and a phase offset. This has
 been successfully applied by Chambers et al. (2014); Uhlemann et al. (2017);
 Holmes et al. (2020). We obtained the parameters of the temperature model
 using field data retrieved from temperature sensors set at 40 and 85 *cm* depth
 280 at the studied site and meteorological data recorded within a year. Resistivity
 values can then be corrected for the effect of the temperature considering :

$$\rho_{corrected} = \rho(1 + \alpha(T_{ref} - T_{model})) \quad (3)$$

where $\rho_{corrected}$ is the resistivity corrected from the effect of the temperature,
 ρ is the initial considered resistivity point, T_{ref} is the reference temperature
 and T_{model} is the estimated temperature at the time and depth of the resistivity
 285 point. We set α for resistivity and normalized chargeability to 0.2 and 0.21

according to the values commonly used in the bibliography and obtained by laboratory experiments (Uhlemann et al., 2017; Coperey et al., 2019; Holmes et al., 2020).

4.4. Extract information from the decay curves through a Debye Decomposition

290 Classical inversion of TDIP used the DC resistivity and the integral chargeability (the partial chargeability integrated over all the time gates) to obtain resistivity and chargeability models (Oldenburg and Li, 1994; Revil et al., 2020). However, no information about the spectral content can be extracted this way. The Debye Decomposition (DD) approach has been used in frequency domain
 295 by Nordsiek and Weller (2008) and in time domain by Tong et al. (2006a); Tarasov and Titov (2007); Martin et al. (2021). In this paper, we follow the approach proposed by Tarasov and Titov (2007) and more recently by Martin et al. (2021) in time domain. The decay curves are converted into an equivalent sum of Debye models via a Relaxation Time Distribution (RTD). The relax-
 300 ation times are logarithmly discretized with 200 steps in a predefined matrix from $10^{-3.5}$ et $10^{3.5}$ s. The *RMS* and the χ^2 values are extracted to ensure the inversion validity. We also define the mean relaxation time τ_{mean} as Nordsiek and Weller (2008) :

$$\tau_{mean} = \exp\left(\frac{\sum_k (m_k \cdot \ln(\tau_k))}{\sum_k m_k}\right) \quad (4)$$

4.5. Clustering

305 Data interpretation is often biased especially as for the quantification of what we consider as a significant variation. Recently in near surface geophysics applications, unsupervised classification has proven to be an efficient tool to define groups of common geophysical characteristics (Ronczka et al., 2017; Whiteley et al., 2021; Carrier et al., 2022) or to identify time-lapse variations (Xu et al.,

2017; Delforge et al., 2021). We apply unsupervised classification to the variations of the different geophysical properties in order to create groups areas where geophysical properties evolve through time in a similar manner. For our specific case, we choose to apply a Gaussian Mixture Model (GMM) for its ability to identify convex clusters in the parameter space and for its speed and easiness of implementation via the scikit-learn python package (Pedregosa et al., 2011). The GMM clustering algorithm is performed to build consistent groups of [resistivity-normalized chargeability-relaxation time] variations. The data space consists of the 70% of data with the highest coverage values, we fixed the number of groups to 5 considering the geological context, the rationale and a BIC (Bayesian information criterion) approach.

4.6. *Surface displacements from ground based images (TSM)*

Thanks to the recent development of high resolution sensors, the use of optical cameras to track surface displacements via image correlations appears to be relevant for landslide monitoring (Gance et al., 2014; Travelletti et al., 2012). These methods have the advantage of being low cost and providing information within the whole image and not at specific points. We used the TSM (Tracing Surface Motion) processing toolbox developed by Desrues et al. (2019) and the VSC (Variable Sequential Correlation) methodology to compute the deformation fields. It consists of a 6 steps methodology based on the use of a time serie of images in a JPG format originated from a monoscopic camera. The 6 steps are detailed in Desrues et al. (2019) and include evaluating the sensor orientation, selecting images, correcting the sensor mouvements, cross-correlating images to get ground motion, extracting the meaningful patterns and applying a geometrical correction.

335 **5. Results**

Time-lapse ERT and IP inversion exhibits satisfying results, RMS ranges from 4.6 to 8.5% and 9.6 to 12.3% for ERT and IP inversion respectively. The resulting cross-plots are shown on figure 2. The residuals for IP inversion are more spread than for ERT inversion which is consistent with the fact that IP data are more affected by noise than ERT data. For the DD, the mean and median values of the χ^2 are 12 and 5 for the September 2021 data set and 5 and 3 for the May 2022 data set respectively which is satisfying. In the framework of the MGL, we do not expect any metallic particles within the materials and we consider that the normalized chargeability is a proxy of surface conductivity and gives information about the amount of fine particles (Mao et al., 2016). The resistivity models exhibit high resistivities in the first 5 m depth (> 4000 Ohm.m). Resistivity is gradually decreasing until 30 m depth to reach values between 200 and 500 Ohm.m. Low values are reached quicker in the downstream part of the MGL (Figure 3, A), D), G)). Normalized chargeabilities values range between 0.001 and 0.3 mS/m. Normalized chargeabilities are less than 0.01 mS/m in the first 5 m depth and then increases with depth to reach 0.3 mS/m in the downstream part of the landslide in May 2022 (Figure 3, B), E), H)). Relaxation time constant model exhibits values in the order of 1 s in the first 15 m depth. The values tend to decrease with depth in September 2021 and increase with depth in May 2022 in the downstream part of the landslide (Figure 3, F), I)). We extracted 3 main hydrogeological units (HU) on the basis of inclinometer data, piezometric measurements, resistivity and normalized chargeability models (figure 3, C). The first one (HU 0, green) consists in the first 5 meters depth of the landslide made up of very loose material, it exhibits very high resistivities, low chargeabilities and variable relaxation time constants. Between 5 and 15 m depth, a second body (HU 1, grey) is defined where resis-

tivities and chargeabilities are intermediate. This agrees with quite remobilized materials, moderately altered micaschists. At depth more than 15 *m*, materials with low resistivity and high chargeability are observed and constitute the third body (HU 2, purple). The latter suggests material with important surface conductivities and agrees with less damaged micaschists. Resistivity variations are represented along the corresponding surface displacements obtained via the TSM method over the considered period of time (figure 4). The variations are defined as $\frac{variable-variable_{ref}}{variable_{ref}}$. Normalized chargeability and mean relaxation time variations are only considered between September 2021 and May 2022 when acquisition parameters and acquisition material are comparable (Figures 5 and 6). The time lapse inversion observations are the following :

- Resistivity tends to globally decrease within the landslide between April and September 2021. The surface displacements associated remains less than 0.1 *m* (figure 4).
- Between September 2021 and May 2022, the resistivity tends to globally increase in the first 15 *m* depth and decreases beyond 15 *m* depth in the downstream part of the landslide. The surface displacements associated are in the order of 0.3-0.4 *m* in the center/center-east-downstream part of the landslide. These results are in agreement with field observations which suggest that the landslide is more active in winter and spring (figure 4). Normalized chargeability generally decreases between September and May for depths more than 15 *m* (figure 5), the decrease being more important in the center/center-east-downstream part of the landslide (up to 80%) than in the uphill part (less than 20% and more often less than 10%). In the subsurface, it generally increases or stays rather stable in the upstream part of the landslide. The mean relaxation time is relatively stable in the uphill part of the landslide. In the downstream part, it slightly decreases

at the outer rings but it exhibits a significant increase of more than 200%
390 in the center/center-east-downstream part of the landslide (figures 5 and
6).

- Between April 2021 and May 2022, resistivity has increased in the center/-
center east-downstream part of the landslide. The surface displacements
associated are the order of 0.3-0.4 *m* in the eastern part of the landslide
395 and can reach 0.8 *m* in the center/center-east-downstream part of the
landslide. Resistivity variations are globally less than 20 % in absolute
value (figure 4).

In the HU 1, resistivity and mean relaxation time tend to increase (from 0
to 20 % and 0 to 50 % respectively) and normalized chargeability tends to
400 decrease (10 to 50 %) between September 2021 and May 2022. Yearly resistivity
variations remain quasi-nul. The variations observed in the HU 0 are more
chaotic. Hence, in the first 15 *m* of the landslide, the variations are season-
sensitive and variations from one year to the other remain quite small. The
HU 2 exhibits few variations in the upstream part of the landslide (less than 10
405 %) whereas a resistivity and normalized chargeability decrease (up to 25 and
50 % respectively) and a strong mean relaxation time increase (more than 200
%) between September 2021 and May 2022 are observed in the downstream
part (figures 5 and 6). Between the end of the summer and the spring, the
piezometer level usually changes from about 30 *m* to 15 *m*, meaning that this
410 deeper geological layer of the landslide fills up with water during the winter
and the spring and drains during the summer. In this hydrogeological unit
we observe seasonal variations but also a non-negligible permanent resistivity
variation in the center/center-east-downstream part. This suggests that some
water could be stored for a longer time at the downstream of the landslide.
415 To help the reading of the variations results, a GMM clustering is applied to

the variations of resistivity, normalized chargeability and mean relaxation time (figure 5 D, E, F). Five groups are extracted, the groups 0 and 1 correspond to groups where resistivity and normalized chargeability have decreased and mean relaxation time has significantly increased between September 2021 and
420 May 2022. The group 2 exhibits low variations for the 3 variables, and the groups 3 and 4 a resistivity decrease, a normalized chargeability increase and a decrease of the mean relaxation time. Groups 0 and 1 are situated directly above the surface areas where surface displacements are the strongest between September 2021 and May 2022. We suggest that the groups 3 and 4 correspond
425 to highly altered material very sensitive to seasonal variations and extreme events (fast snow melting, heavy rains) and represent $17500 m^3$. The groups 0 and 1 corresponds to materials subject to seasonal and permanent variations, the latter controlling the overall large scale dynamics of the landslide. Groups 0 and 1 correspond to a volume of 6000 and 11000 m^3 respectively.

430 6. Discussion

The measured TDIP data appears to be of good quality and consistent from one campaign to the other. The resulting normalized chargeabilities and mean relaxation times and their variations are spatially consistent which suggests that the results are trustworthy. However, we recorded field decay curves only on
435 a limited time range (0.05 to 8 s). It implies a known information loss at early times (i.e. highest frequencies) and later times (i.e. lowest frequencies). Even if some useful information did come out from the time range used here, we suggest that more information about the spectral content of the IP response could be extracted with the use of a more elaborated post-processing technique as the
440 one describe by Olsson et al. (2016). This work has not been performed here and could be the subject of the further developement of a freely available module. It

has been revealed that the pulse length influences the magnitude and the shape of the decay curves (Mao et al., 2016; Olsson et al., 2019), so we choose to work with a fixed pulse length. We did not study the influence of changing the pulse
445 length on our data mainly because of time issues. Considering these limitations, we interpreted the RTD and extracted information (mean relaxation time) carefully with full knowledge and only qualitatively.

The nature and origin of resistivity seasonal variations within landslides are strongly site and geology dependent (Supper et al., 2014; Gance et al., 2016;
450 Palis et al., 2017b,a; Uhlemann et al., 2017; Holmes et al., 2020). What's more, it appears that the resistivity variations observed in clayey landslides are quite tenuous, generally less than 10% (Gance et al., 2016). Between September and May at the MGL, we observed either a resistivity increase or decrease whether we consider surface or deep landslide material respectively. We assume from
455 field observations and piezometric levels that the water saturation increases from summer to winter within the landslide even if we can not quantify it precisely. To illustrate and better understand the resistivity variations through the seasons in the different parts of the MGL, we computed the theoretical evolution of the total conductivity as a function of water saturation at different
460 water conductivities (i.e. different salinities) for a typical clay/silt material taking surface conductivity into account following Revil et al. (2020) model. The figure 7 presents the evolution of bulk conductivity with saturation for pore filled of snow melt water (0.2 mS/m), rain water (2 mS/m) or ground water (about 20 mS/m with almost no seasonal variations, publicly available
465 water quality measurements for the different sources in the Arly valley). It illustrates the fact that when water saturation increases within the landslide during winter and spring, the conductivity can decrease within the subsurface material due to the very low conductivity of the brought fluids (green arrow on

the figure 7) whereas, meanwhile, it can increase in the deeper parts of the
470 landslide due to an increase of water content of constant medium salinity.

The relation between normalized chargeability and saturation or relaxation
time and saturation is strongly material dependant. As a consequence, the po-
larization models inferred to explain the results are not universal (Titov et al.,
2004; Kemna et al., 2005; Binley et al., 2005; Ghorbani et al., 2009; Revil, 2012;
475 Breede et al., 2012; Okay et al., 2013; Abdulsamad et al., 2017; Zhang et al.,
2019; Martin et al., 2022). In our case, the total chargeability of our mixture of
damaged/crushed micascists and silts never exceed 40 mV/V and no metallic
particles are observed which suggests that our media might not act as pure clay
neither like material mixed with metallic particles. Its behaviour can then be
480 compared to the one in sands and/or mixtures of sand and clay. The polariza-
tion mechanisms can then be explained by a combination of the short narrow
pore model (Titov et al., 2004) and a possible rearrangement of clay particles in
the pore space (Breede et al., 2012). In this type of medium, the relaxation time
constant increases with saturation and normalized chargeability increases and
485 then decreases after reaching a certain saturation threshold. Salinity may also
affect polarization mechanisms (Revil, 2012; Lévy et al., 2019b; Martin et al.,
2022), its influence also strongly depends on the dominant polarization and
conduction mechanisms. Experiments realized on sands, sandstones and clay
agreed on the fact that polarization tend to increase with salinity until reaching
490 a threshold and decrease at high salinities (Revil and Skold, 2011; Weller and
Slater, 2012; Weller et al., 2013; Hördt et al., 2016; Mendieta et al., 2021). In
the HU 0 (surface unit), the normalized chargeability tend to globally increase
between September 2021 and May 2022 in parallel to a salinity decrease and
a saturation increase from a very low starting saturation. The previous exper-
495 iments would expect a normalized chargeability increase with saturation and

a decrease with salinity. As a consequence, we assume here that the saturation effect is dominant on the observed response. In the HU 1, the normalized chargeability decreases when the soil get saturated with low conductivity fluids. This could be the result of the combined influence of water saturation increase
500 from a medium to high starting saturation and a salinity decrease. In the HU 2, we assumed the salinity remains constant (on the basis of the measured salinities in the sources around) so the observed decrease of normalized chargeability and the increase of relaxation times highlight the increase in water saturation. Considering all the previous observations, we infer that the main sources of feeding
505 water in the subsurface units (HUs 0 and 1) are the rain and snow melt events whereas in the HU 2 (deeper unit) the water circulating within the massif is the main source of additional fluids.

IP measurements are related to the surface area of the pore network as a consequence they are related to the pore size and can help predicting the permeability (Börner et al., 1996; Slater, 2007; Revil, 2012; Weller et al., 2015). Recently, several studies proposed the used of laboratory calibrated relationships between permeability and IP parameters to estimate the permeability at the field scale (Fiandaca et al., 2018; Maurya et al., 2018; Revil et al., 2020). In all the studies, the estimation of the permeability implies to consider some fixed parameters (like porosity or water conductivity) and/or some assumption about petrophysical models. Some studies also proposed to link the permeability with the relaxation time but the relationships appear to be strongly dependant of the geology (Binley et al., 2005; Tong et al., 2006b; Zisser et al., 2010). Considering these limitations, we decided to apply the relationship proposed by Weller et al. (2015) and applied on the field by Orozco et al. (2022) to infer an order of magnitude of the permeability and its evolution through time using the normalized chargeability as a proxy of the imaginary conductivity in order

to avoid formulation implying the use of the formation factor which we can not properly spatially define in the case of MGL. Hydraulic conductivity is then defined as :

$$K = \frac{g \cdot d_w}{\mu} \frac{3.47 \cdot 10^{-16} \sigma_{DC}^{1.11}}{\sigma''^{2.41}} \quad (5)$$

where g is the gravitational constant, d_w and μ are groundwater density and dynamic viscosity in $kg \cdot m^{-3}$ and $kg \cdot m^{-1} s^{-1}$ respectively. The relationship
510 between the normalized chargeability and the quadrature conductivity is related to the number of decades of the frequency range considered such as $Mn \approx \alpha \sigma''$ with $\alpha \approx \frac{2}{\pi} \ln(A)$, where $A = 10^k$, k being the number of decades between minimum and maximum acquisition frequencies. In our case $\alpha \approx 4$. Hydraulic conductivity decreases from $10^{-5} m/s$ in the subsurface of the MGL to 10^{-8}
515 m/s at 30 m depth within the landslide (Figure 9). After heavy rains or quick snow melt episodes resurgences appear within a few days at the base of the landslide, suggesting that the loose surface material are permeable. Considering a 5 m depth layer, the vertical hydraulic conductivity would be the order of $10^{-5} m/s$. The saturation and the resulting pore pressure increase of these
520 loose subsurface materials generates rock flows. These rock flows happened to be almost permanent during the rainy season. Figure 8 represents the rain and cumulative rain measured at the site along with the surface displacements during the year 2021. The cross-correlation between the signals exhibits a maximum at a lag of 25 to 30 days for the targets 4 to 15 situated in the active part of
525 the landslide. Considering a 15 m depth sliding surface, the resulting vertical hydraulic conductivity would be the order of $10^{-7} m/s$. We underline the fact that the orders of magnitude of hydraulic conductivity estimated from the field observations are consistent and validate the estimation made from geophysical data.

530 The MGL results from the toppling failure of the sub-vertical micaschists.

The latter are strongly foliated rocks composed of sheets of phyllosilicates. In such materials, hydraulic conductivity is strongly dependant on the direction and so is polarization (Kelly and Reiter, 1984; Cosenza et al., 2007; Börner et al., 2018; Abdulsamad et al., 2019). In the case of the MGL, the horizontal
535 bending of the micaschists generated by the toppling failure would generate low vertical hydraulic conductivities in the 15 to 30 *m* depth unit. But this same layer could be more permeable in the opposite direction which would allow the fluids from within the massif to circulate. This layer (15 to 30 *m* depth) is acting as an aquifer and undergo seasonal water table variations (Figure 10). During
540 winter and spring water infiltrates from the upper massif and gather in this layer at the downstream of the landslide. As a consequence, the pore pressure increases and the available space for electrical charges conduction widens which lead to a decrease of normalized chargeability and resistivity, an increase in mean relaxation time and an increase in permeability. At this point, general
545 mass movements of the landslide appear. In winter and spring, when HU 2 is pressurized the units above are more prone to destabilization and respond quickly to rain and snow melt events. We suggest that, the observed surface displacements results from the combination of rain/snow melt infiltration and underground water accumulation at depth.

550 7. Conclusion

This paper presents the results of the first ERT and TDIP time-lapse inversion performed on an active landslide. A methodology developed to extract the spectral information from IP data through a DD has been applied in order to extract as many information as possible from the TDIP data. The resistivity
555 variations are quite small within the landslide and are not sufficient to extract significant information about the unstable volumes. It appears that IP data are

decisive for the estimation of these volumes. The results of the time-lapse ERT and TDIP are very well correlated with the surface displacements obtained via the TSM method. The use of a non-supervised classification method enables to build groups on the basis of the temporal evolution of the variables without 560 formulating any hypothesis about the physical mechanisms or using petrophysical models. The time-lapse inversion of the geophysical data and the clustering enables us to discriminate between weather sensitive subsurface volume with potentially fast reaction times and deeper volume affected by seasonal water 565 level variation but also exhibiting permanent variations. The first represents at least 17500 m^3 of materials and the latter situated above the highest surface displacements represents a total of 17000 m^3 potentially unstable materials. These results are critical for the establishment and sizing of security systems nearby the landslides. TDIP data have the potential to provide additional information 570 about the water content and the geometry of fluid flow circulations within landslides as long as we understand how the considered variables evolves with water content and/or salinity. Hence, TDIP appears to be a very promising tool in landslide characterization and dynamics assessment. Note that, the possibilities could be enhanced by the improvement of the data processing methodology and 575 the use of complementary geophysical data such as refraction seismic data. This paper did not focus on the theory of the IP, neither on the establishment of a petrophysical model for the studied site. In order to better understand the physical processes at stake, especially the polarization mechanisms at work from a polarization theory point of view, more laboratory data and study would be 580 needed.

Declarations

7.1. Consent for publication

The authors have the consent for publication of the data.

7.2. Conflict of interest/Competing interests

585 The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Acknowledgments

This work has been funded by the POIA-FEDER european program in the
590 framework of the MIROIR project (PA0020589). We would like to thank the BRGM and especially Edouard Equilbey for supporting the MIROIR project. We are also grateful to the CD73 and especially Anne Lescurier for providing data and feedbacks about the observed dynamics of the landslide. We would like to thank Nicolas Moreau, Charles Ribeyre and Mathilde Masson for their
595 precious help on the field. We also warmly thank Thomas Gunther and Tina Martin for their very useful advices. We also thank the two anonymous reviewers for their constructive comments which greatly helped to improve the quality of the manuscript.

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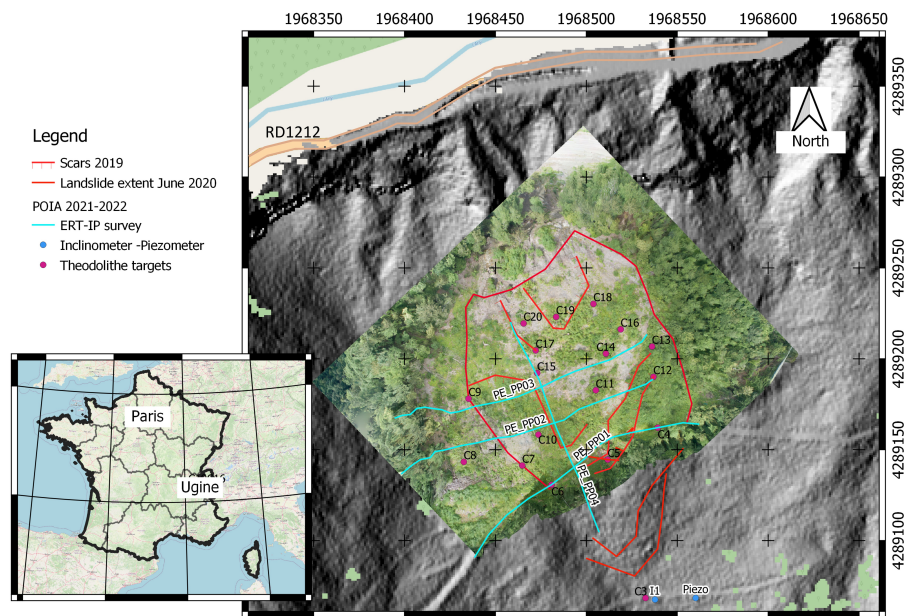


Figure 1: Spatial location of the studied area. The spatial extend of the most active part of the landslide is represented by red lines. The location of the geophysical surveys (blue lines), inclinometers-piezometer (blue points) and theodolithe targets (pink points) are also represented.

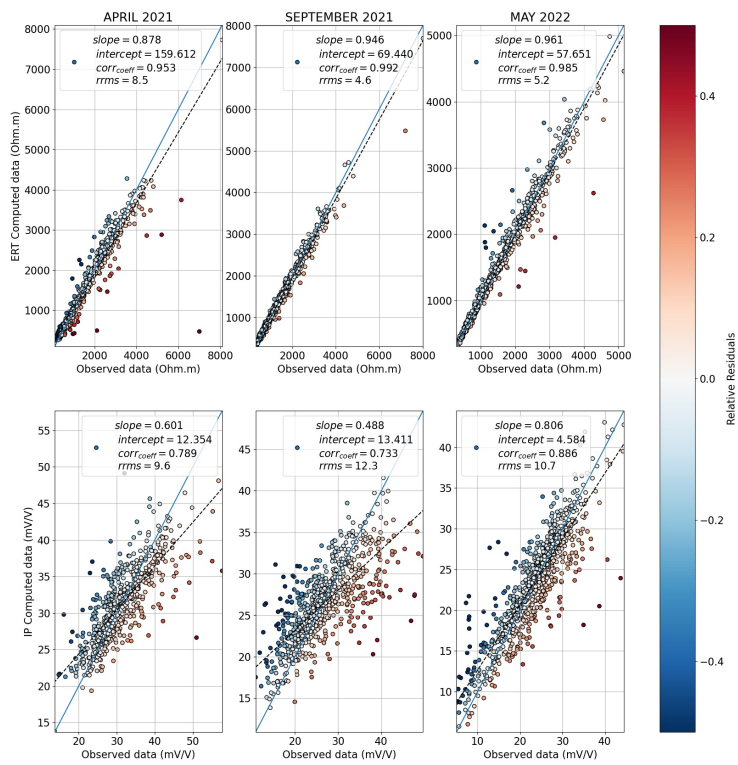


Figure 2: Cross-plots obtained from the ERT and IP time-lapse inversion. The color scales with the relative residuals.

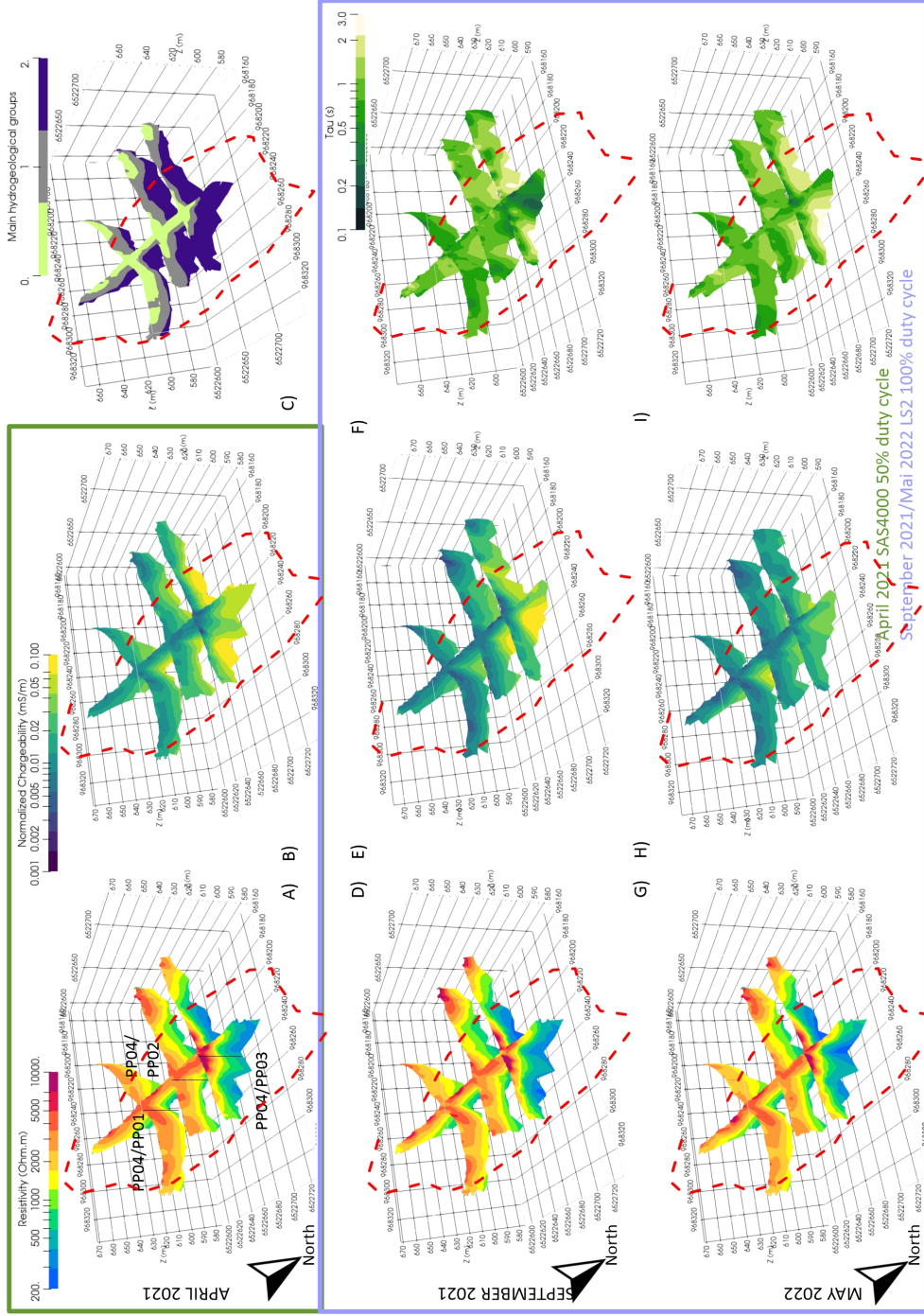


Figure 3: Resistivity (A, D, G), Normalized chargeability (B, E, H) and mean relaxation time (F, I) models obtained from the time-lapse inversion for the April 2021, September 2021 and May 2022 campaigns. C) Main Hydrogeological units defined on the basis of inclinometer data, piezometric measurements and geophysical models. Black lines represent the location of the further extracted vertical profiles.

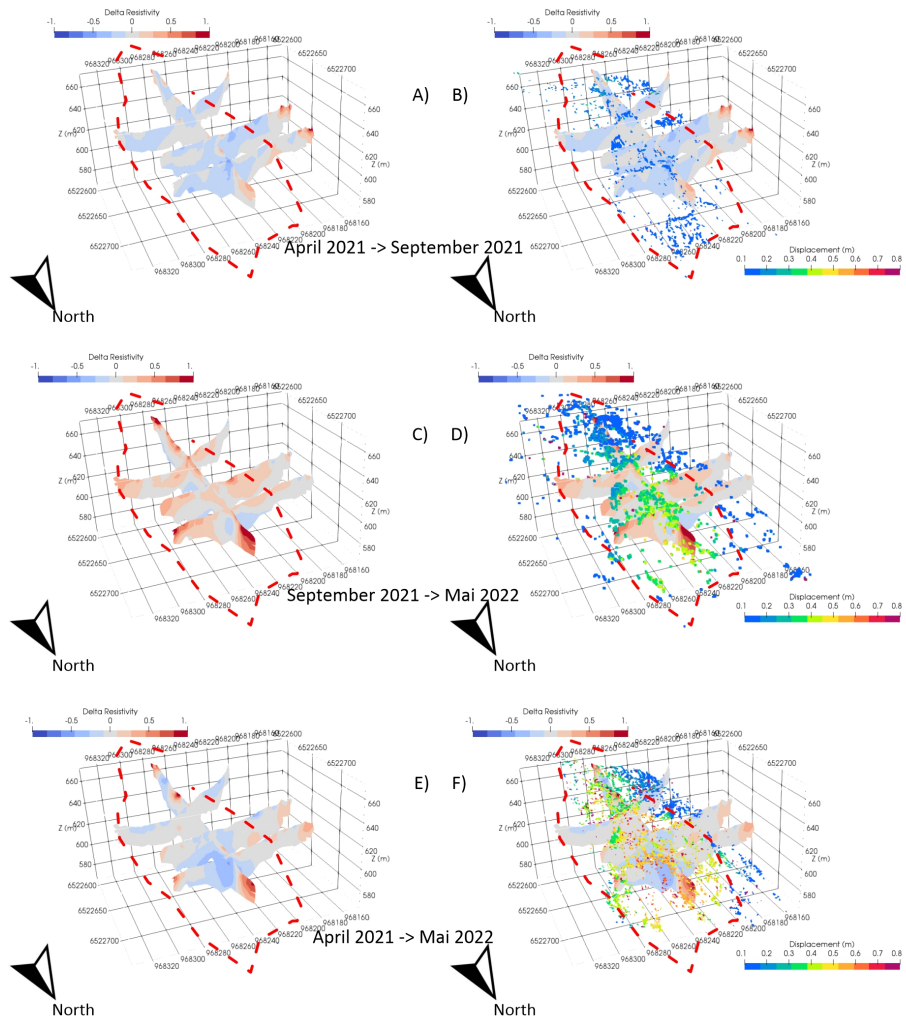


Figure 4: Relative resistivity variations (A, C, E) and their corresponding surface displacement obtained via TSM methodology (B, D, F). Relative variations are expressed in % and displacements in meters.

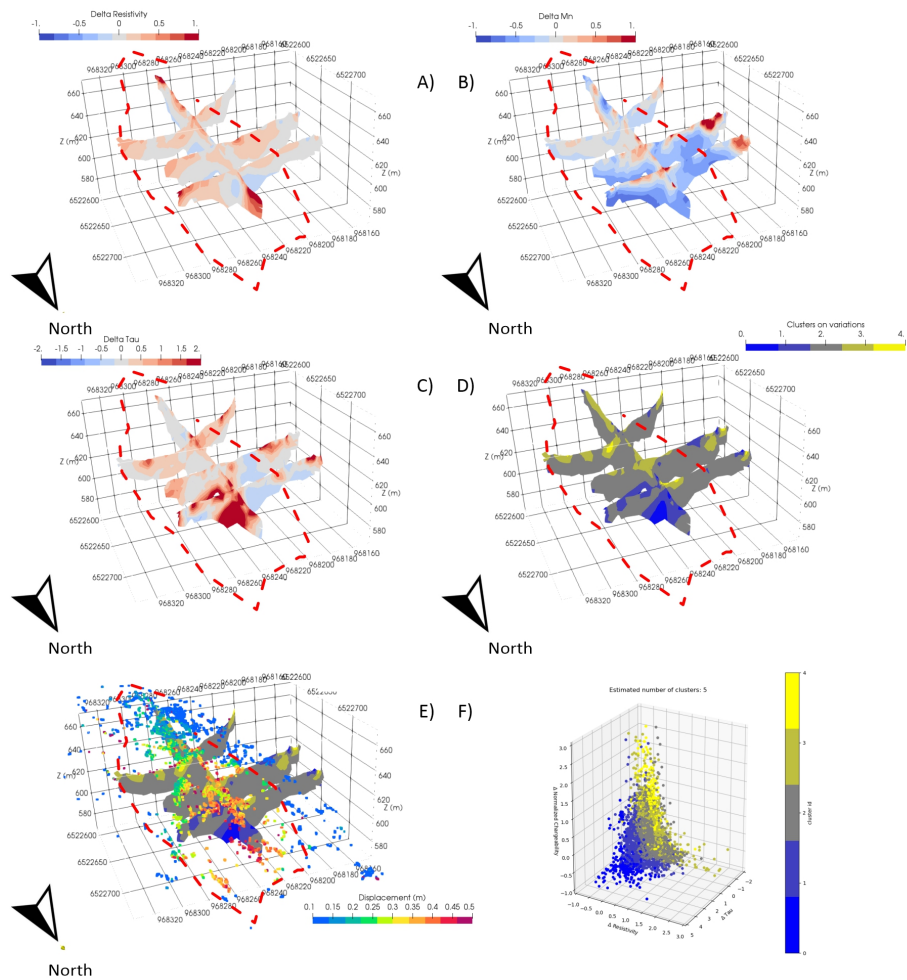


Figure 5: Relative resistivity (A), normalized chargeability (B), mean relaxation time (C) variations between September 2021 and May 2022. The results of the clustering applied on relative variations and its confrontation to surface displacements is presented on panels D, E and F. The relative variations are defined as $\frac{variable - variable_{ref}}{variable_{ref}}$, where september 2021 is used as reference.

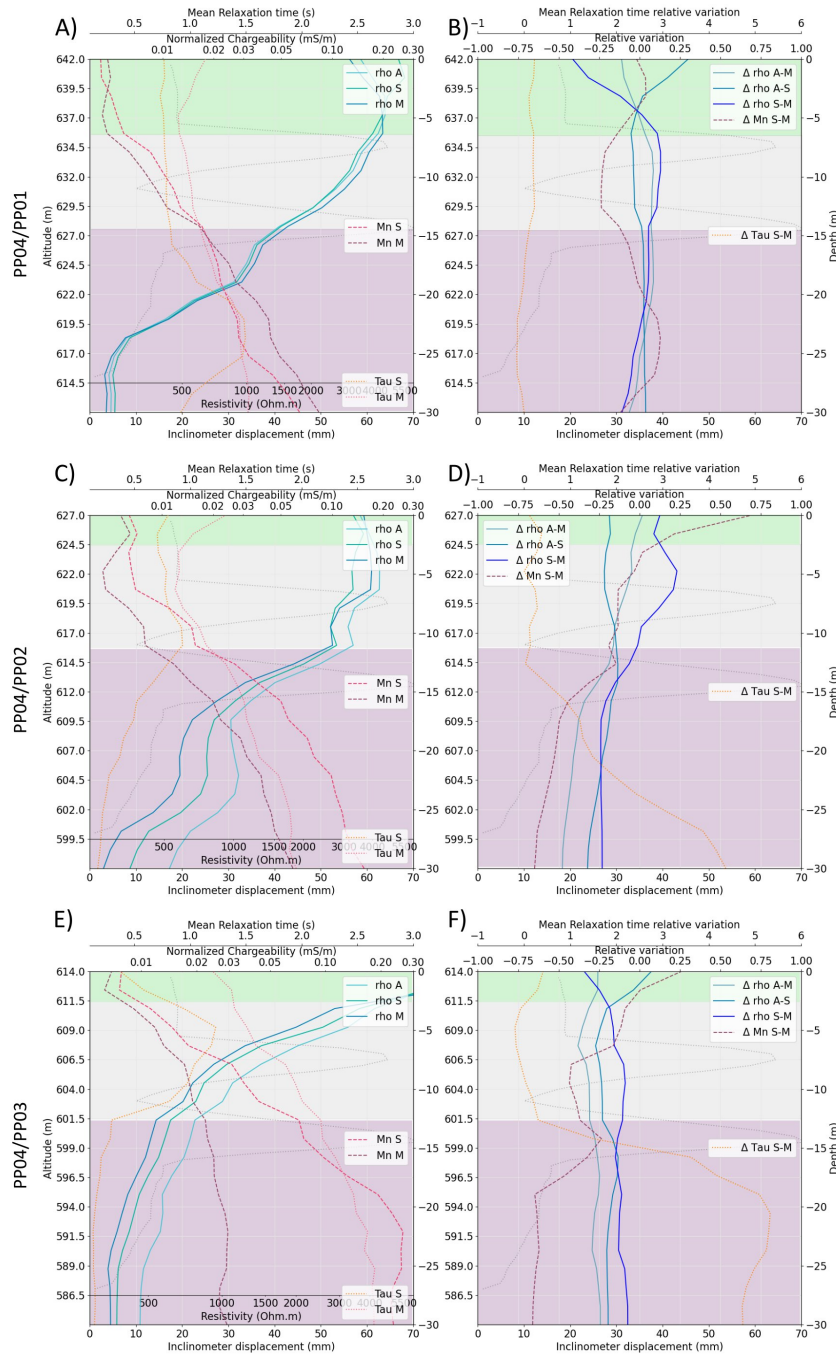


Figure 6: Resistivity (ρ in $Ohm.m$), normalized chargeability (Mn in mS/m) and mean relaxation times (Tau in s) at the three intersections between PP04 and PP01 (A), PP02 (C) and PP03 (E) respectively from top to bottom of the landslide, location on figure 3. Relative resistivity, normalized chargeability and mean relaxation time variations at the three intersections between PP04 and PP01 (B), PP02 (D) and PP03 (F) respectively from top to bottom of the landslide (relative variations are defined as $\frac{variable - variable_{ref}}{variable_{ref}}$ and have therefore no units). Please note that relative variations scale is different for B panel compared to D and F to ensure readability. The inclinometer measured displacements in mm are shown in grey in the background. The letters A, S and M in the captions stand for April 2021, September 2021 and May 2022. The green, gray and purple backgrounds correspond to the three main hydrogeological units within the landslide based on the interpretation of the inclinometer, piezometric and geophysical data.

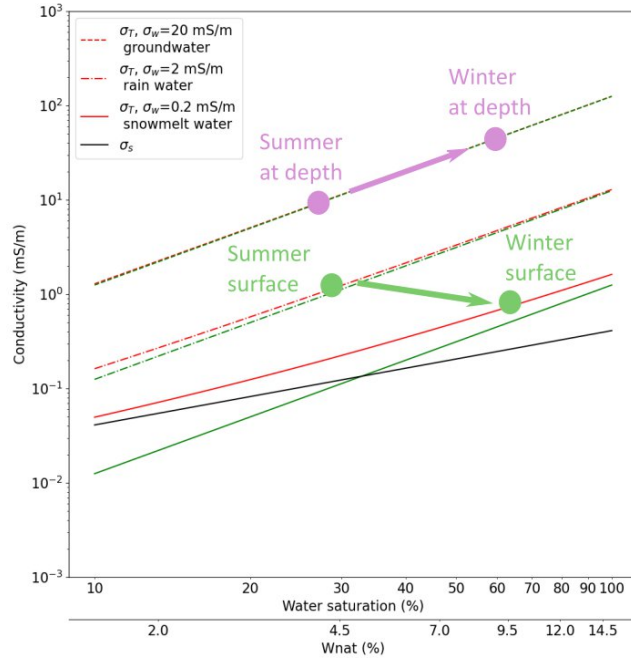


Figure 7: Theoretical influence of water saturation on total, hydraulic and surface conductivity for different water conductivities for a typical material made of silt or clay silt. Red lines represent the total (bulk) DC conductivity (σ_T), green lines represent the electrolytic conductivity (σ_H) and black line represents surface conductivity (σ_s). σ_w stands for water conductivity. Green and purple arrows represent the possible evolution at the surface and at 20 m depth within the MGL from the end of the summer to the spring. We consider the total DC conductivity as : $\sigma_T = \sigma_H + \sigma_s = \theta^m \sigma_w + \theta^{m-1} \rho_g (B - \lambda) CEC$, where θ is the volumetric water content, m the saturation exponent further set to 2, ρ_g the grain density ($kg.m^{-3}$), B and λ the apparent mobility of counterions for surface conduction and polarization respectively and set to 3.1×10^{-9} and $3 \times 10^{-10} m^2.s^{-1}.V^{-1}$ and CEC the cation exchange capacity in $meq/100g$ further set to 20 (Revil et al., 2020).

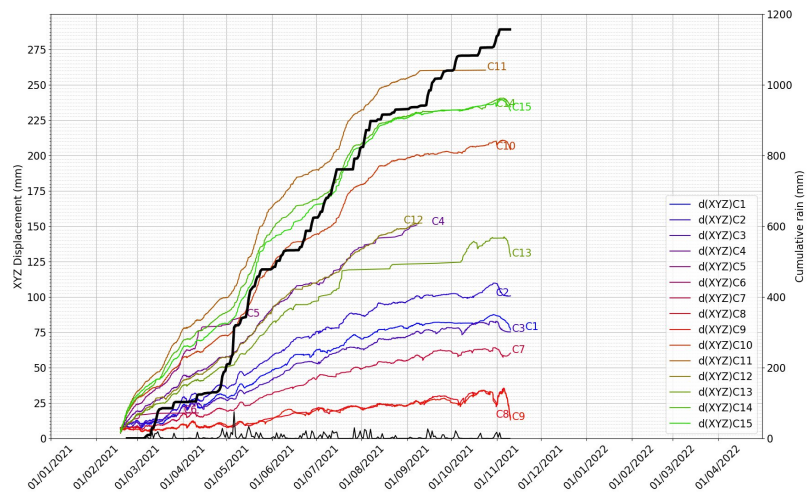


Figure 8: Surface displacements measured from February to November 2021 using an automated theodolite. Black curves represent rain (thin black line) and cumulated rain (bold black line) in *mm*. Orange points represent the shearing measured at -15 m at the inclinometer I1 ahead of the upper end of PP04 profile in *mm*.

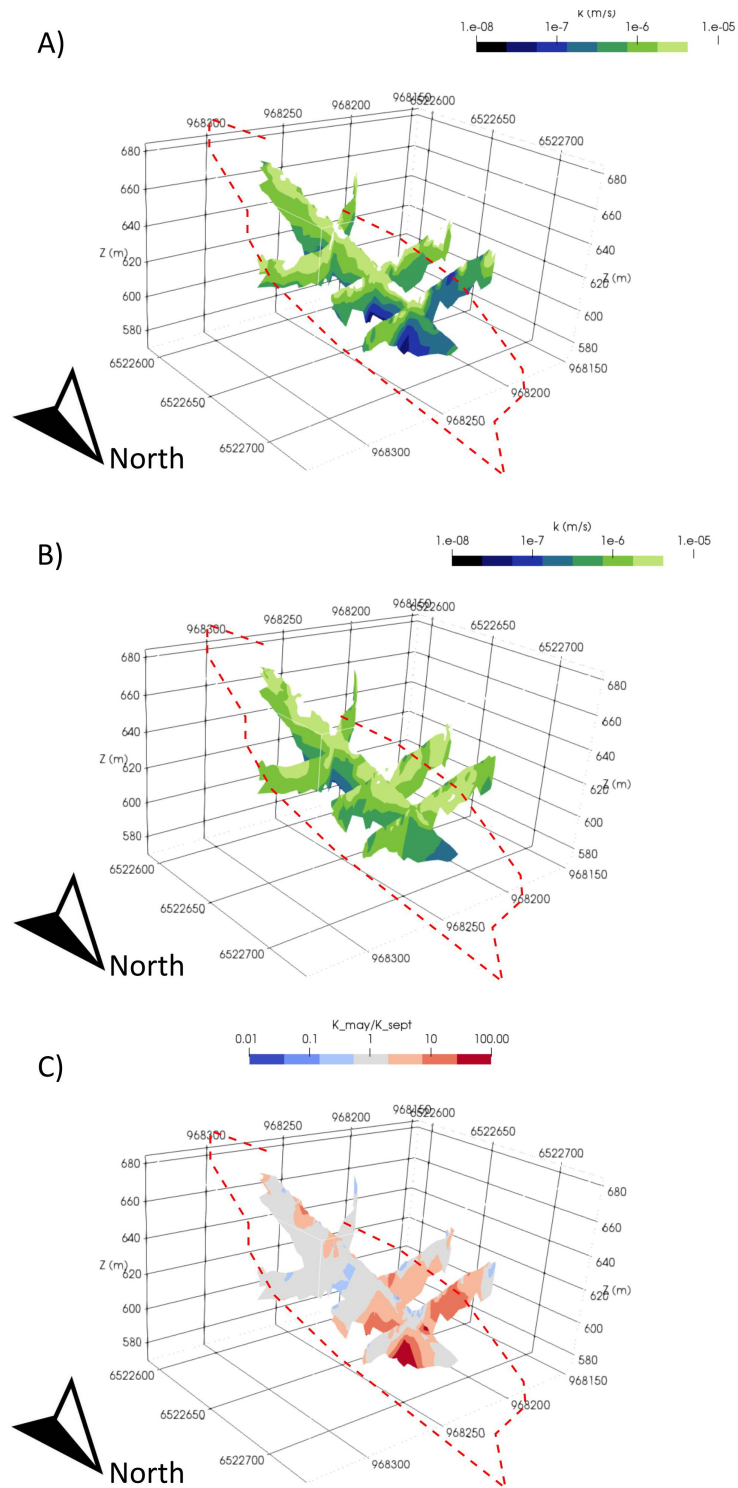


Figure 9: Permeability models for September 2021 (A), May (2022) and their ratio (C) extracted from Weller et al. (2015) relationship.

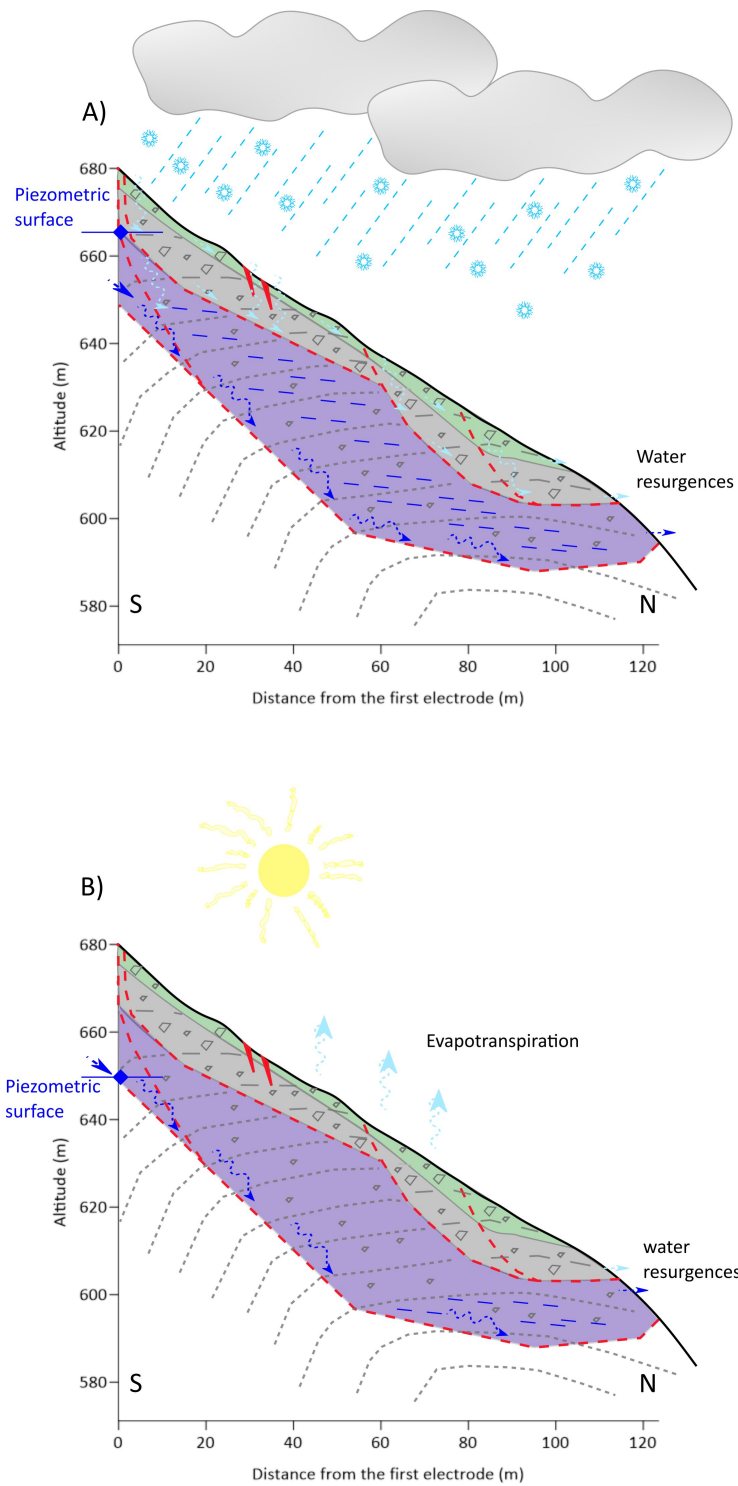


Figure 10: Simplified cross-section of the MGL landslide along the PP04 profile showing subsurface hydrogeological fluid dynamics through the seasons (A) December to May B) June to November). The three main bodies (green, grey, purple) are the three main hydrogeological bodies identified in the paragraph 5. The damage degree of the micaschist is increasing with depth. Toppling is beginning at the interface between the bedrock and the purple body and damage progressively increases toward the surface in the two other units. Red dashed lines represent the landslide surfaces inferred from inclinometer data and geophysical data. The purple body acts as a semi-permeable aquifer unit and piezometric surface varies within this body. Soil moisture accumulation is possible at the bottom of the landslide. Red triangles represent open cracks act as permeable barriers.